

Boundary value problems for harmonic functions on domains in Sierpinski gaskets

Hua Qiu

(Joint work with [Shiping Cao](#))

Nanjing University

Email: huaqiu@nju.edu.cn

June 17, 2017, [Cornell](#)

Dirichlet problem

For a domain Ω , finding a function which is *harmonic* in Ω that takes prescribed values on $\partial\Omega$,

$$\begin{cases} \Delta u = 0 \text{ in } \Omega, \\ u|_{\partial\Omega} = f, f \in C(\partial\Omega). \end{cases} \quad (1)$$

Dirichlet problem

For a domain Ω , finding a function which is *harmonic* in Ω that takes prescribed values on $\partial\Omega$,

$$\begin{cases} \Delta u = 0 \text{ in } \Omega, \\ u|_{\partial\Omega} = f, f \in C(\partial\Omega). \end{cases} \quad (1)$$

For a bounded domain Ω with *sufficiently smooth boundary* $\partial\Omega$, the problem is always *solvable*,

$$u(x) = \int_{\partial\Omega} f(s) \partial_n G(x, s) ds, \quad \forall x \in \Omega.$$

Dirichlet problem

For a domain Ω , finding a function which is *harmonic* in Ω that takes prescribed values on $\partial\Omega$,

$$\begin{cases} \Delta u = 0 \text{ in } \Omega, \\ u|_{\partial\Omega} = f, f \in C(\partial\Omega). \end{cases} \quad (1)$$

For a bounded domain Ω with *sufficiently smooth boundary* $\partial\Omega$, the problem is always *solvable*,

$$u(x) = \int_{\partial\Omega} f(s) \partial_n G(x, s) ds, \quad \forall x \in \Omega.$$

$G(x, y)$ is the *Green's function*, $\partial_n G(x, s)$ is the *Poisson kernel* for Ω .

Harmonic extension algorithm (w.r.t. the standard energy)

- Sierpinski gasket \mathcal{SG} the “ $\frac{1}{5} - \frac{2}{5}$ ” rule
- \mathcal{SG}_3 the “ $\frac{8}{15} - \frac{4}{15} - \frac{3}{15}$ ” rule
- $\mathcal{SG}_l, l \geq 2, \mathcal{SG}_l = \bigcup_{i=0}^{\frac{l^2+l-2}{2}} F_i \mathcal{SG}_l$
 $V_0 = \{q_0, q_1, q_2\}, V_m = \bigcup_{i=0}^{\frac{l^2+l-2}{2}} F_i V_{m-1}$

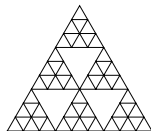
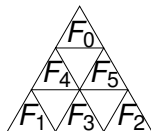
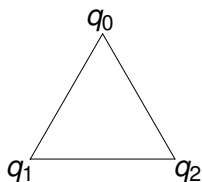
Solving linear equations

$$\sum_{y \sim_1 x} (h(y) - h(x)) = 0, \forall x \in V_1 \setminus V_0$$

harmonic extension matrices A_i 's: $h \circ F_i|_{V_0} = A_i h|_{V_0}$

local algorithm: $h \circ F_{wi}|_{V_0} = A_i h \circ F_w|_{V_0}$

- p.c.f. self-similar set (omit)



The problem

Consider Dirichlet problems on *connected subsets* of p.c.f. self-similar sets (*with fractal boundary*)

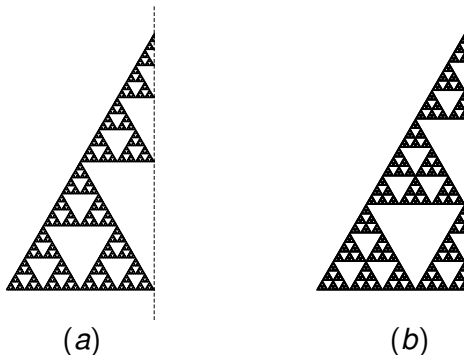
we will focus on subdomains in $\mathcal{SG}_l, l \geq 2$

The problem

Consider Dirichlet problems on *connected subsets* of p.c.f. self-similar sets (*with fractal boundary*)

we will focus on subdomains in \mathcal{SG}_l , $l \geq 2$

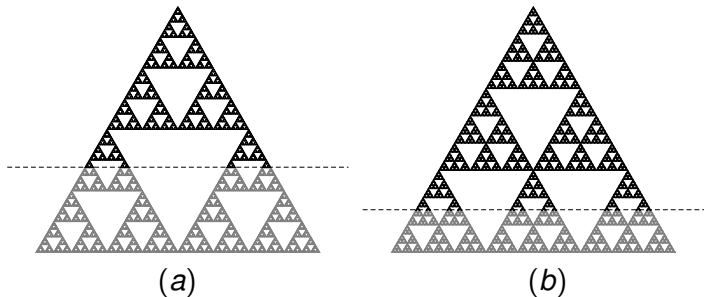
- half domains



Half domains of \mathcal{SG} and \mathcal{SG}_3 .

The problem

- upper domains (lower domains)



Upper and lower domains in \mathcal{SG} and \mathcal{SG}_3 .

Proposition 1.

Let Ω be a *half, upper or lower* domain in $S\mathcal{G}_l$, $l \geq 2$. The Dirichlet problem

$$\begin{cases} \Delta u = 0 \text{ in } \Omega, \\ u|_{\partial\Omega} = f, f \in C(\partial\Omega) \end{cases}$$

has a unique solution.

Our goal:

The problem

Proposition 1.

Let Ω be a *half, upper or lower* domain in $S\mathcal{G}_l$, $l \geq 2$. The Dirichlet problem

$$\begin{cases} \Delta u = 0 \text{ in } \Omega, \\ u|_{\partial\Omega} = f, f \in C(\partial\Omega) \end{cases}$$

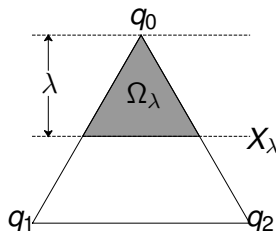
has a unique solution.

Our goal:

- to find the exact harmonic extension algorithm in Ω
- to estimate the energy of harmonic solutions in terms of boundary values

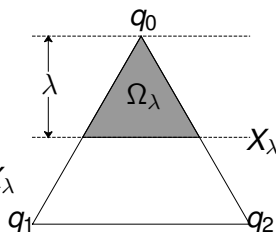
Upper domains in \mathcal{SG} (Strichartz, 1999; Owen and Strichartz, 2012; Guo, Kogan, Q. and Strichartz, 2014)

- Extension algorithm
- Energy estimate



Upper domains in \mathcal{SG} (Strichartz, 1999; Owen and Strichartz, 2012; Guo, Kogan, Q. and Strichartz, 2014)

- Extension algorithm
- Energy estimate



Main tool: Haar series expansion along X_λ
Haar basis $\{\phi_w\} \cup \{1\}$ for $L^2(X_\lambda, \mu)$
the standard measure μ on X_λ

The only **generator** of the Haar basis is **antisymmetric**, its extended harmonic solution is also **antisymmetric**, so can be **localized** to any small scales along the boundary X_λ .

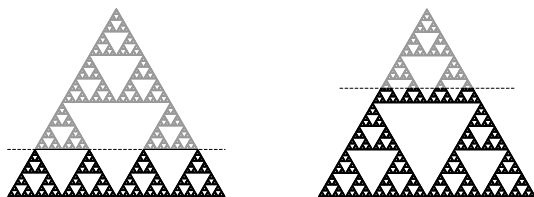
Need symmetry, could not be extended to $\mathcal{SG}_l, l \geq 3$

Prior works

Lower domains in \mathcal{SG} (Strichartz, 1999)

very special case, the domains are made up of 2^m adjacent triangles of size 2^{-m} lying on the bottom line of \mathcal{SG}

- *the boundary is a finite set*
- *the energy estimate is unknown*



Two typical domain Ω_λ^- 's, λ is (or not) a dyadic rational.

For general lower domains, we have little knowledge

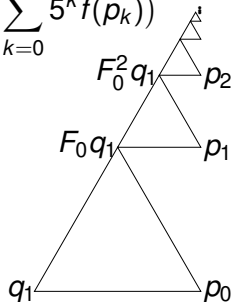
Half domain of \mathcal{SG} (Li and Strichartz, 2014)

- Extension algorithm

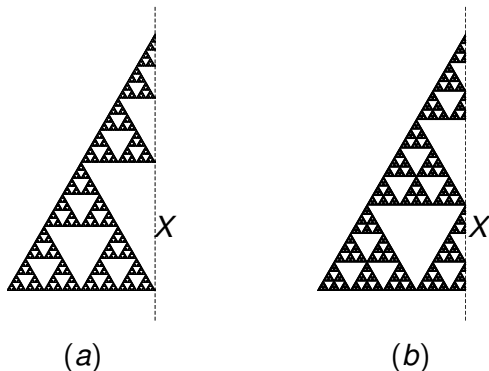
$$u(F_0^m q_1) = \frac{1}{5^m} (f(q_1) - \frac{3}{7} \sum_{k=0}^{\infty} (\frac{1}{3})^k f(p_k) + \frac{10}{7} \sum_{k=0}^{\infty} 5^k f(p_k)) + \frac{3}{7} \sum_{k=0}^{\infty} (\frac{1}{3})^k f(p_{m+k})$$

- Energy estimate

$$\mathcal{E}(u) \asymp (f(q_1) - f(p_0))^2 + \sum_{k=0}^{\infty} (\frac{5}{3})^k (f(p_k) - f(p_{k+1}))^2$$



- Method: solving systems of infinite linear equations, multiplying infinite matrices



Half domains of S_G and S_{G_3} .

For S_G , X consists a countable infinite set

For S_{G_l} , $l \geq 3$, X is a Cantor set

The approach is not applicable for S_{G_l} , $l \geq 3$

- There exist *crucial points* for certain domains. We only need to solve finite linear equations.
- The normal derivative of the harmonic solution at a boundary point is determined by the integral of boundary data w.r.t. a signed measure.

Extension 1: Half domain of \mathcal{SG}_3

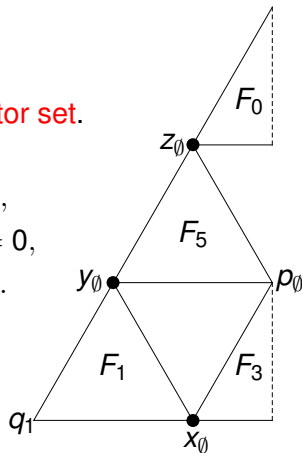
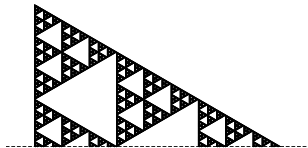
- Structure of half domain

$$\bar{\Omega} = F_1 \mathcal{SG}_3 \cup F_5 \mathcal{SG}_3 \cup F_0 \bar{\Omega} \cup F_3 \bar{\Omega}.$$

The boundary set $X = \partial\Omega \setminus \{q_1\}$ is a **Cantor set**.

- Basic equations

$$\begin{cases} \frac{7}{15} \partial_n^{\leftarrow} u(x_\emptyset) + 2u(x_\emptyset) - u(y_\emptyset) - f(q_1) = 0, \\ 4u(y_\emptyset) - u(x_\emptyset) - u(z_\emptyset) - f(q_1) - f(p_\emptyset) = 0, \\ \frac{7}{15} \partial_n^{\leftarrow} u(z_\emptyset) + 2u(z_\emptyset) - u(y_\emptyset) - f(p_\emptyset) = 0. \end{cases}$$



Extension 1: Half domain of \mathcal{SG}_3

- Calculation of $\partial_n^{\leftarrow} u(x_\emptyset)$ and $\partial_n^{\leftarrow} u(z_\emptyset)$ in terms of boundary data

Observation 1.

$$\partial_n^{\leftarrow} u(x_\emptyset) = \frac{15}{7} \partial_n^{\leftarrow} (u \circ F_3)(q_1),$$

$$\partial_n^{\leftarrow} u(z_\emptyset) = \frac{15}{7} \partial_n^{\leftarrow} (u \circ F_0)(q_1).$$

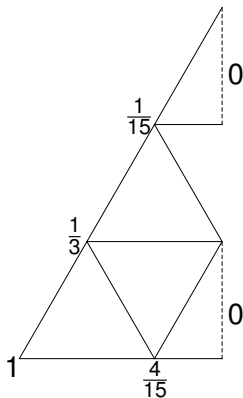
Observation 2.—Gauss-Green's formula

Roughly, we have

$$\mathcal{E}_\Omega(u, h_a) = \partial_n^{\leftarrow} u(q_1)$$

$$= \int_{\partial\Omega} u(x) \partial_n h_a(x) d\mu(x),$$

h_a being the restriction of a **antisymmetric harmonic function** on Ω .

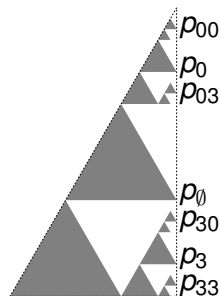
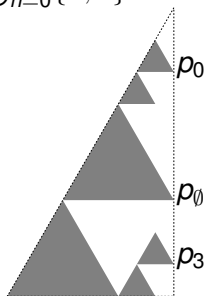
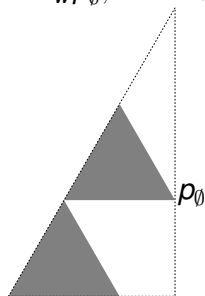


Extension 1: Half domain of \mathcal{SG}_3

By using a sequence of *simple sets* to approximate $\bar{\Omega}$, we prove

$$\partial_n^{\leftarrow} u(q_1) = 3f(q_1) + \sum_{w \in \tilde{W}_*} \partial_n^{\rightarrow} h_a(p_w) f(p_w).$$

$$p_w = F_w p_\emptyset, w \in \tilde{W}_* = \bigcup_{n=0}^{\infty} \{0, 3\}^n.$$



Extension 1: Half domain of \mathcal{SG}_3

Introduce the (probability) measure

$$\mu_a = \frac{1}{3} \sum_{w \in \tilde{W}_*} (-\partial_n^{\rightarrow} h_a(p_w)) \delta_{p_w},$$

then

$$\partial_n^{\leftarrow} u(q_1) = 3f(q_1) - 3 \int_X f d\mu_a (= \mathcal{E}_\Omega(u, h_a) \text{ if } u \in \text{dom} \mathcal{E}_\Omega).$$

For this reason, we may view the measure $3\delta_{q_1} - 3\mu_a$ as the normal derivative of h_a along $\partial\Omega$.

Exact data: $-\partial_n^{\rightarrow} h_a(p_w) = \frac{6}{7} \prod_{i=1}^{|w|} \mu_{w_i}$, for $\mu_0 = \frac{1}{7}, \mu_3 = \frac{4}{7}$.

Theorem 1.1. Extension algorithm

There exists a unique solution of the Dirichlet problem (1) on the half domain of SG_3 . In addition, we have

$$u(x_\emptyset) = \frac{4}{15}f(q_1) + \frac{1}{15}f(p_\emptyset) + \frac{1}{30} \int_X f \circ F_0 d\mu_a + \frac{19}{30} \int_X f \circ F_3 d\mu_a,$$

$$u(y_\emptyset) = \frac{1}{3}f(q_1) + \frac{1}{3}f(p_\emptyset) + \frac{1}{6} \int_X f \circ F_0 d\mu_a + \frac{1}{6} \int_X f \circ F_3 d\mu_a,$$

$$u(z_\emptyset) = \frac{1}{15}f(q_1) + \frac{4}{15}f(p_\emptyset) + \frac{19}{30} \int_X f \circ F_0 d\mu_a + \frac{1}{30} \int_X f \circ F_3 d\mu_a.$$

Theorem 1.2. Energy estimate

We have the energy estimate that

$$C_1 Q(f) \leq \mathcal{E}(u) \leq C_2 Q(f),$$

$$Q(f) = (f(q_1) - f(p_\emptyset))^2$$

where

$$+ \sum_{w \in \tilde{W}_*} \left(\frac{15}{7}\right)^{|w|} \left((f(p_w) - f(p_{w0}))^2 + (f(p_w) - f(p_{w3}))^2 \right).$$

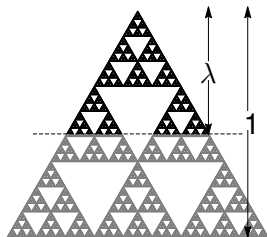
Remark 1. Theorem 1.2 is motivated by Li and Strichartz, 2014.

Remark 2. Both Theorem 1.1 and 1.2 can be extended to \mathcal{SG}_l cases.

Extension 2: Upper domains of \mathcal{SG}_3

- **Notations.** Assume $q_0 = (\frac{1}{\sqrt{3}}, 1)$, $q_1 = (0, 0)$, $q_2 = (\frac{2}{\sqrt{3}}, 0)$. For $0 < \lambda \leq 1$, define the *upper domain*

$$\Omega_\lambda = \{(x, y) \in \mathcal{SG}_3 \setminus V_0 \mid y > 1 - \lambda\},$$



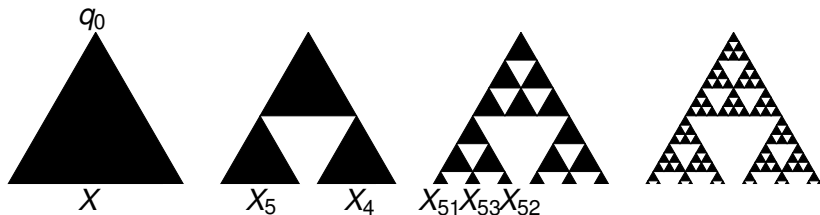
Extension 2: Upper domains of $\mathcal{S}\mathcal{G}_3$

Write $\lambda = \sum_{k=1}^{\infty} \iota_k 3^{-m_k}$, with $0 < m_1 < m_2 < \dots$ and $\iota_k = 1$ or 2 .
The boundary is $\partial\Omega = \{q_0\} \cup X_\lambda$, with X_λ homeomorphic to

$$\Sigma_\lambda = \prod_{k=1}^{\infty} S_{\iota_k},$$

where $S_1 = \{4, 5\}$, $S_2 = \{1, 2, 3\}$.

Denote the *cylinder* of X corresponding to $w \in W_m^\lambda = \prod_{k=1}^m S_{\iota_k}$ by X_w .

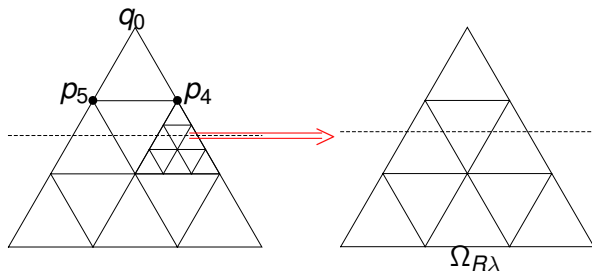


Extension 2: Upper domains of \mathcal{SG}_3

- Basic equations

Case 1. ($\iota_1 = 1$) $\lambda = (1 + R\lambda)3^{-m_1}$, with $0 < R\lambda \leq 1$

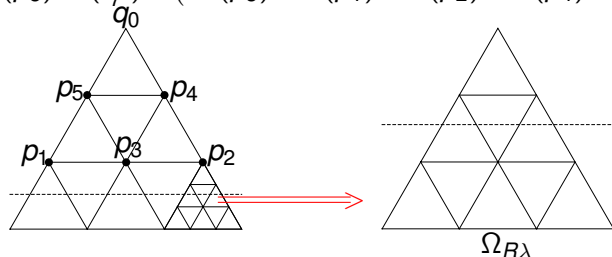
$$\begin{cases} \partial_n^\uparrow u(p_5) + \left(\frac{15}{7}\right)^{m_1} (2u(p_5) - u(p_4) - f(q_0)) = 0, \\ \partial_n^\uparrow u(p_4) + \left(\frac{15}{7}\right)^{m_1} (2u(p_4) - u(p_5) - f(q_0)) = 0. \end{cases}$$



Extension 2: Upper domains of \mathcal{SG}_3

Case 2. ($\nu_1 = 2$) $\lambda = (2 + R\lambda) \cdot 3^{-m_1}$, with $0 < R\lambda \leq 1$

$$\begin{cases} 4u(p_5) - u(p_1) - u(p_3) - u(p_4) - f(q_0) = 0, \\ 4u(p_4) - u(p_2) - u(p_3) - u(p_5) - f(q_0) = 0, \\ \partial_n^\uparrow u(p_1) + \left(\frac{15}{7}\right)^{m_1} (2u(p_1) - u(p_3) - u(p_5)) = 0, \\ \partial_n^\uparrow u(p_2) + \left(\frac{15}{7}\right)^{m_1} (2u(p_2) - u(p_3) - u(p_4)) = 0, \\ \partial_n^\uparrow u(p_3) + \left(\frac{15}{7}\right)^{m_1} (4u(p_3) - u(p_1) - u(p_2) - u(p_4) - u(p_5)) = 0. \end{cases}$$



Extension 2: Upper domains of \mathcal{SG}_3

- Calculation of Normal derivatives

Still the Gauss-Green's formula

$$\partial_n^\uparrow u(q_0) = \partial_n^\uparrow h_0(q_0) f(q_0) - \int_{X_\lambda} f d\mu_\lambda (= \mathcal{E}_{\Omega_\lambda}(u, h_0) \text{ if } u \in \text{dom} \mathcal{E}_{\Omega_\lambda}). \quad (2)$$

h_0 is harmonic in Ω_λ with $h_0(q_0) = 1$ and $h_0 \equiv 0$ on X_λ .

Observation 1. $-\mu_\lambda$ is the “normal derivative” of h_0 on X_λ

Observation 2. $\partial_n^\uparrow h_0(p_i) = (\frac{15}{7})^{m_1} h_0(p_i) \partial_n^\uparrow h_0^{R\lambda}(q_0)$

Calculate μ_λ on each cylinder X_w inductively using Observation 2.

The proof for formula (2) involves a series of *simple sets* approximating to the domain.

Extension 2: Upper domains of \mathcal{SG}_3

The **extension algorithm** is derived by solving the basic equations.

Theorem 2.1. Extension algorithm

*There exists a unique solution of the Dirichlet problem (1) on the upper domain Ω_λ . In addition, we have **explicit extension algorithm** for $u(p_i)$'s in terms of $f(q_0)$ and integrals of boundary value data w.r.t. $\mu_{R\lambda}$.*

Extension 2: Upper domains of \mathcal{SG}_3

- Haar series

Definition:

If $\iota_{n+1} = 1$, define $\psi_w^{(1)} = \chi_{X_{w4}} - \chi_{X_{w5}}$, for $w \in W_n^\lambda$.

If $\iota_{n+1} = 2$, define

$$\begin{cases} \psi_w^{(1)} = \chi_{X_{w1}} - \chi_{X_{w2}} \\ \psi_w^{(2)} = \mu_w^{-1} (\mu_{w3} \chi_{X_{w1}} - 2\mu_{w1} \chi_{X_{w2}} + \mu_{w3} \chi_{X_{w2}}) \end{cases}, \text{ for } w \in W_n^\lambda.$$

$$\begin{array}{ccc} \psi_\emptyset^{(1)} & \underline{1} & \underline{-1} \\ \psi_\emptyset^{(2)} & \underline{\mu_3} & \underline{-2\mu_1} \end{array} \quad \begin{array}{ccc} \psi_\emptyset^{(1)} & \underline{1} & \underline{0} \\ \psi_\emptyset^{(2)} & \underline{\mu_3} & \underline{-2\mu_1} \end{array} \quad \begin{array}{ccc} \psi_\emptyset^{(1)} & \underline{0} & \underline{-1} \\ \psi_\emptyset^{(2)} & \underline{-2\mu_1} & \underline{\mu_3 = \mu_2} \end{array}$$

Properties:

- $\{\psi_w^{(j)}\}_{w \in W_n^\lambda} \cup \{1\}$ form an **orthogonal basis** in $L^2(X_\lambda, \mu_\lambda)$.
- $\|\psi_w^{(j)}\|_{L^2(X_\lambda, \mu_\lambda)} \asymp \mu_w^{\frac{1}{2}}$

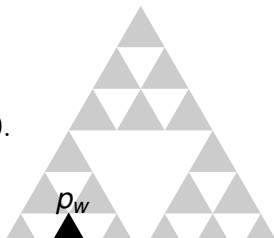
Extension 2: Upper domains of \mathcal{SG}_3

Let $h_w^{(j)}$ be the harmonic functions with

$$h_w^{(j)}|_{X_\lambda} = \psi_w^{(j)}, \quad h_w^{(j)}(q_0) = 0.$$

Then we have

1. $h_w^{(j)}$ supports locally, $\partial_n^\uparrow h_w^{(j)}(p_w) = 0$
2. $\mathcal{E}_\Omega(h_w^{(j)}) \asymp (\frac{15}{7})^{m|w|+1}$.
3. $\{h_w^{(j)}\} \cup \{h_0\}$ are **pairwisely orthogonal in energy**.



Extension 2: Upper domains of $\mathcal{S}\mathcal{G}_3$

A brief explanation for the **orthogonality in energy**:

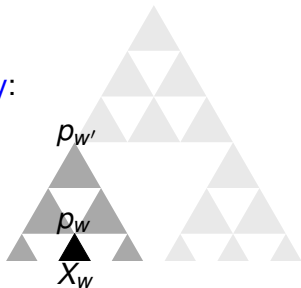
Case 1. $X_w \cap X_{w'} = \emptyset$: Disjoint support.

Case 2. $w = w', j \neq j'$: Different symmetry.

Case 3. $X_w \subset X_{w'}$:

$$\begin{aligned}\mathcal{E}_{\Omega_\lambda}(h_w^{(j)}, h_{w'}^{(j')}) &= \mathcal{E}_{\text{supp}h_w^{(j)}}(h_w^{(j)}, h_{w'}^{(j')}) \\ &= \partial_n^\uparrow h_w^{(j)}(p_w) h_{w'}^{(j')}(p_w) + \int_{X_w} h_{w'}^{(j')} \partial_n h_w^{(j)} \\ &= 0 + \int_{X_w} C \partial_n h_w^{(j)} = 0\end{aligned}$$

(The total of normal derivative on X_w is zero.)



Extension 2: Upper domains of \mathcal{SG}_3

Theorem 2.2. Energy Estimate

Let u be the harmonic function in Ω_λ with boundary values $u(q_0) = a$ and $u|_{X_\lambda} = f$, where

$$f = b + \sum_{w \in W_*^\lambda} \sum_{j \leq \ell_{|w|+1}} c_w^{(j)} \psi_w^{(j)}.$$

Then $\mathcal{E}_{\Omega_\lambda}(u)$ is bounded above and below by a multiple of

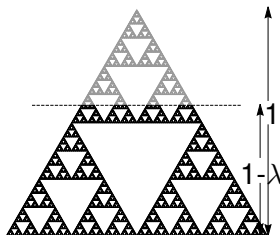
$$\left(\frac{15}{7}\right)^{m_1} (a - b)^2 + \sum_{n=0}^{\infty} \sum_{w \in W_n^\lambda} \sum_{j=1}^{\ell_{n+1}} \left(\frac{15}{7}\right)^{m_{n+1}} |c_w^{(j)}|^2.$$

Remark. Both Theorem 2.1 and 2.2 can be extended to \mathcal{SG}_l cases.

Extension 3: Lower domains of \mathcal{SG}

- **Notations.** Assume $q_0 = (\frac{1}{\sqrt{3}}, 1)$, $q_1 = (0, 0)$, $q_2 = (\frac{2}{\sqrt{3}}, 0)$. For $0 \leq \lambda < 1$, define the lower domain

$$\Omega_\lambda^- = \{(x, y) \in \mathcal{SG} \setminus V_0 \mid y < 1 - \lambda\}$$



Consider the **binary expansion**

$$\lambda = 0.e_1 e_2 \cdots = \sum_{k=1}^{\infty} e_k(\lambda) 2^{-k}.$$

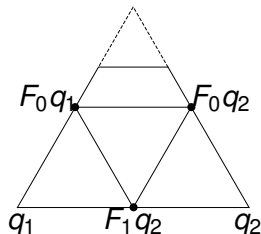
Define $S\lambda = 0.e_2 e_3 \cdots$.

Extension 3: Lower domains of \mathcal{SG}

- Basic equations

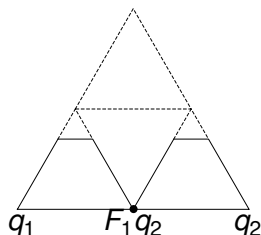
Case 1. $e_1 = 0$

$$\begin{cases} 4u(F_1q_2) - f(q_1) - f(q_2) - u(F_0q_1) - u(F_0q_2) = 0, \\ \frac{3}{5}\partial_n^{\leftarrow} u(F_0q_1) + (2u(F_0q_1) - u(F_1q_2) - f(q_1)) = 0, \\ \frac{3}{5}\partial_n^{\rightarrow} u(F_0q_2) + (2u(F_0q_2) - u(F_1q_2) - f(q_2)) = 0. \end{cases}$$



Case 2. $e_1 = 1$

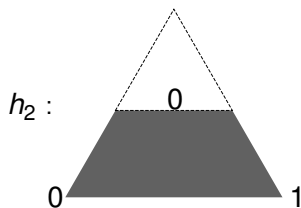
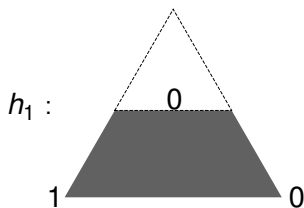
$$\partial_n^{\rightarrow} u(F_1q_2) + \partial_n^{\leftarrow} u(F_1q_2) = 0.$$



Extension 3: Lower domains of \mathcal{SG}

- **Basic ideas**

Introduce two harmonic functions h_1, h_2 .



Calculate $\partial_n h_1, \partial_n h_2$ on the boundary.

Use Gauss-Green's formula to get the normal derivatives, and solve the basic equations.

- We have the **extension algorithm** (omit), could be extended to \mathcal{SG}_I case.

However, the **energy estimate** remains **open**.

- Dirichlet to Neumann map
- trace problem and extension problem
- other Sobolev types of norms
- Green's function
- spectral decimation and spectral analysis
- other subdomains in p.c.f. self-similar sets

Thank you!