About Lie Groups

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Contents

1	Lie	Groups	1
	1.1	Beginning Details	1
	1.2	The Exponential Map and Useful Curves	3
2	The	e Lie Algebra of a Lie Group	4
	2.1	General Lie Algebras	4
	2.2	The Tangent Space at the Identity	4
	2.3	Left Invariant Vector Fields	5
	2.4	$T_e G \cong \mathcal{L}(G)$ as Vector Spaces	7
	2.5	$T_eG \cong \mathcal{L}(G)$ as Lie Algebras	7

1 Lie Groups

1.1 Beginning Details

A Lie group is a group G with the structure of a smooth manifold, such that the two maps

$$G \to G, x \mapsto x^{-1}$$
 and $G \times G \to G, (x, y) \mapsto xy$

are smooth. This last condition is equivalent to requiring that the single map $G \times G \to G, (x, y) \mapsto xy^{-1}$ be smooth.

A Lie group G comes with a lot of structure. There is a distinguished element $e \in G$, the group identity. This means there is a distinguished tangent space, $T_eG \subset TG$. For each $g \in G$, we obtain three maps:

- (1) left multiplication: $L_q: G \to G, h \mapsto gh;$
- (1) right multiplication: $R_q: G \to G, h \mapsto hg$; and
- (1) conjugation: $\Psi_g \colon G \to G, \ h \mapsto L_g \circ R_{g^{-1}}(h) = R_{q^{-1}} \circ L_g(h) = ghg^{-1}.$

Each of these maps is smooth, and in fact they are all diffeomorphisms. The inverses of L_g , R_g , and Ψ_g are $L_{g^{-1}}$, $R_{g^{-1}}$, and $\Psi_{g^{-1}}$, respectively. They are also group homomorphisms, so they are **Lie group isomorphisms**. Also, note that the left and right multiplication maps commute with each other. For all $g, h \in G$, we have $L_g \circ R_h = R_h \circ L_h$.

Remark 1.1. For whatever reason, most of Lie theory is centered around the left multiplication maps, but it could just as well have been developed using the right multiplication maps.

The three maps above are *canonical* with respect to the Lie group structure. Therefore all tangent spaces of G are canonically isomorphic. For $g, h \in G$, we have the canonical linear isomorphisms

$$T_g(L_{hg^{-1}}): T_gG \to T_hG \qquad and \qquad T_g(L_{gh^{-1}}): T_hG \to T_gG.$$

Thus all tangent spaces of G are canonically isomorphic to the distinguished tangent space of G, T_eG .

Let $g \in G$. Because $\Psi_g(e) = geg^{-1} = e$, we have a canonical operator on the distinguished tangent space T_eG , given by $T_e\Psi_g \colon T_eG \to T_eG$. We denote this map by $\operatorname{Ad}(g)$. Since Ψ_g is a diffeomorphism, $\operatorname{Ad}(g) = T_e\Psi_g$ is a linear isomorphism, so $\operatorname{Ad}(g) \in \mathfrak{GL}(T_eG)$, so

$$\operatorname{Ad}: G \to \mathfrak{GL}(T_eG)$$

is a group representation of G, called the **adjoint representation**.

Recall that $\mathfrak{GL}(T_eG)$ is the inverse image of the open set $\mathbb{R} \setminus \{0\}$ under the continuous (and smooth) map det: $\mathfrak{gl}(T_eG) \to \mathbb{R}$, so it is an open subset of the vector space $\mathfrak{gl}T_eG$. Thus $\mathfrak{GL}(T_eG)$ is a smooth manifold with tangent bundle $\mathfrak{GL}(T_eG) \times \mathfrak{gl}(T_eG)$. Therefore the tangent map of Ad at the identity e is a map T_e Ad: $T_eG \to \mathfrak{gl}(T_eG)$. By slightly restructuring the domain and codomain, we obtain a map

ad:
$$T_eG \times T_eG \to T_eG$$
, $(v, w) \mapsto \operatorname{ad}(v)w := (T_e\operatorname{Ad})(v)w$.

Note that ad is linear in both v and w, so ad is bilinear.

1.2 The Exponential Map and Useful Curves

For each Lie group G, we have the **exponential map**, $\exp_G: T_eG \to G$. We usually omit the subscript from exp if there is no confusion. It is defined by means of **one-parameter subgroups**, which we will not discuss here. The exponential map is characterized by the fact that if $v \in T_eG$ and $s, t \in \mathbb{R}$, then

$$\exp\left((s+t)v\right) = \exp(sv) \cdot \exp(tv) \left(=\exp(tv) \cdot \exp(sv)\right),$$

and the following Lemma.

Lemma 1.2. Let $v \in T_eG$, and let $c \colon \mathbb{R} \to G$ be the smooth curve given by $t \mapsto \exp(tv)$. Then $\dot{c}(0) = v$.

Note that $\mathfrak{GL}(T_eG)$ is a Lie group under multiplication, and that its tangent space at the identity is essentially $\mathfrak{gl}(T_eG)$. Therefore we have a map

$$\exp_{\mathfrak{GL}(T_eG)} \colon \mathfrak{gl}(T_eG) \to \mathfrak{GL}(T_eG).$$

Since $\operatorname{ad}(v) \in \mathfrak{gl}(T_eG)$ for all $v \in T_eG$, we have

$$\exp_{\mathfrak{GL}T_eG}\left(\mathrm{ad}(v)\right) \in \mathfrak{GL}(T_eG).$$

It is natural to ask what element of $\mathfrak{GL}(T_eG)$ this might be.

Theorem 1.3. Let $v \in T_eG$. Then

$$Ad(\exp_G v) = \exp_{\mathfrak{GL}(T_eG)} \left(adv \right)$$

Remark 1.4. Dropping the subscripts from the exponential maps, we obtain the commutative diagram



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Combining Theorem 1.3 and Lemma 1.2 yields the following result.

Proposition 1.5. Let $v, w \in T_eG$, and let $c: \mathbb{R} \to T_eG$ be the smooth curve given by $t \mapsto Ad(\exp(tv))(w)$. Then

$$\dot{c}(0) = ([X,Y]).$$

2 The Lie Algebra of a Lie Group

2.1 General Lie Algebras

A Lie algebra is a real vector space L equipped with a skew-symmetric bilinear map $L \times L \to L, (v, w) \mapsto [v, w]$, called a **bracket**, which satisfies the **Jacobi identity**

$$[u, [v, w]] + [v, [w, u]] + [w, [u, v]] = 0$$

for all $u, v, w \in L$. Two standard examples are the set of vector fields on a manifold with the Lie bracket, or the set of $n \times n$ real (or complex) matrices with the bracket

$$[A,B] := AB - BA.$$

2.2 The Tangent Space at the Identity

The tangent space T_eG at the identity is a real vector space. Using the three classes of maps inherent in the Lie group structure, we can equip T_eG with a bracket that makes it a Lie algebra. The vector space T_eG with this bracket is denoted \mathfrak{g} , and called the Lie algebra of the Lie group G.

Each step in the construction of the Lie bracket for \mathfrak{g} is *natural*, in the sense that it is preserved by smooth homomorphisms between Lie groups. Let H be another Lie group and $\rho: G \to H$ be a smooth homomorphism. The naturality of each step below will be shown by a commutative diagram involving G, H, and ρ .

$$\begin{array}{c} G \xrightarrow{\rho} H \\ \Psi_g \middle| & & \downarrow \Psi_{\rho(g)} \\ G \xrightarrow{\rho} H \end{array}$$

As described above, for each $g \in G$, we obtain a Lie group isomorphism $\Psi_g \colon G \to G$ and a linear isomorphism $\operatorname{Ad}(g) \colon T_e G \to T_e G$. Then $\operatorname{Ad}(g) \in \mathfrak{GL}(T_e G)$, so we have a smooth homomorphism $\operatorname{Ad} \colon G \to \mathfrak{GL}(T_e G)$.

$$\begin{array}{c|c} T_eG \xrightarrow{T_e\rho} T_eH \\ Ad(g) & & \downarrow Ad(\rho(g)) \\ T_eG \xrightarrow{T_e\rho} T_eH \end{array}$$

We define a bracket on T_eG by [v, w] = ad(v)w. It remains to be shown that this bracket is anti-symmetric and satisfies the Jacobi identity. We will not prove this here, although it will follow from the fact that the Lie bracket of vector fields satisfies these properties.

$$\begin{array}{c|c} T_eG \xrightarrow{T_e\rho} T_eH \\ ad(v) & \downarrow ad((T_e\rho)v) \\ T_eG \xrightarrow{T_e\rho} T_eH \end{array}$$

2.3 Left Invariant Vector Fields

Definition 2.1. Let $f: \mathcal{M} \to nfold$ be a diffeomorphism between smooth manifolds, and let $X \in \mathfrak{X}(\mathcal{M})$ and $Y \in \mathfrak{X}(\mathcal{N})$. The **pushforward** of X by f is

$$f_*(X) := Tf \circ X \circ f^{-1} \in \mathfrak{X}(\mathcal{N}),$$

and the **pullback** of Y by f is

$$f^*(Y) := Tf^{-1} \circ Y \circ f \in \mathfrak{X}(\mathcal{M}).$$

Note that

$$f^*(Y) = (f^{-1})_*(Y)$$
 and $f_X(X) = (f^{-1})^*(X)$.

Definition 2.2. A vector field $X \in \mathfrak{X}(G)$ is called **left invariant** if

$$(T_h L_g) (X(h)) = X (L_g(h)) = X(gh)$$

for all $g, h \in G$. This means the following diagram commutes for each $g \in G$.

$$\begin{array}{c} G \xrightarrow{X} TG \\ L_g \downarrow & \downarrow^{T_e L_g} \\ G \xrightarrow{X} TG \end{array}$$

The set of all left invariant vector fields on G is denoted $\mathcal{L}(G)$.

Remark 2.3. Let $X \in \mathcal{L}(G)$. Then $TL_g \circ X = X \circ L_g$ for all $g \in G$. Thus

$$TL_g \circ X \circ (L_g)^{-1} = X$$
 and $(TL_g)^{-1} \circ X \circ L_g = X$

for all $g \in G$. Certainly if a vector field satisfies either of the above equations for all $g \in G$ it must be left invariant. Therefore $\mathcal{L}(G)$ is the set of vector fields invariant under pushforward by left multiplication, which is also the set of vector fields invariant under pullback by left vector fields.

The set $\mathcal{L}(G)$ is clearly a real vector space, but it is not clear what its dimension is. There's no reason to assume that the dimension be finite, but it is. It's actually quite a surprise.

Recall the Lie bracket of vector fields. This can be defined in terms of flows of vector fields, or in terms of derivations. Let $X, Y \in \mathfrak{X}(G)$ be vector fields, let Φ_X^t, Φ_Y^t denote their respective flows, and let $\mathscr{D}_X, \mathscr{D}_Y$ denote their respective associated derivations. Then the **Lie bracket** $[X, Y] \in \mathfrak{X}(G)$ is the unique vector field such that

$$[X,Y] = \mathscr{L}_X Y := \frac{d}{dt} \big|_{t=0} (\Phi_X^t)^* Y,$$

or equivalently,

$$\mathscr{D}_{[X,Y]} = \mathscr{D}_X \circ \mathscr{D}_Y - \mathscr{D}_Y \circ \mathscr{D}_X$$

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The pushforward of vector fields by diffeomorphisms preserves the Lie bracket [see page 144 in Conlon's *Differentiable Manifolds, 2nd Edition*]. Since left invariant vector fields can be categorized as those that are invariant under pushforward by all left multiplications, this implies that the Lie bracket of two left invariant vector fields is also left invariant. Therefore $\mathcal{L}(G)$ equipped with the Lie bracket is a Lie alegra.

2.4 $T_e G \cong \mathcal{L}(G)$ as Vector Spaces

We have two Lie algebras associated with G: the tangent space at the identity, T_eG , with the bracket induced by ad, and the left invariant vector fields, mathcalL(G), with the Lie bracket. In this section we will demonstrate that they are isomorphic as vector spaces.

Define a map $\nu \colon T_e G \to \mathfrak{X}(G)$ by

$$\nu_{\xi}(g) = T_e L_g(\xi)$$

for all $\xi \in T_e G$ and $g \in G$. Because tangent maps are linear, so is ν . For all $\xi \in T_e G$ and $g, h \in G$ we have

$$(T_h L_g) \big(\nu_{\xi}(h) \big) = (T_h L_g) \big(T_e L_h(\xi) = T_e (L_g \circ L_h)(\xi) = T_e L_{gh}(\xi) = \nu_{\xi}(gh) = (\nu_{\xi} \circ L_g)(h).$$

Therefore ν_{ξ} is left invariant, so ν really is a map $T_e G \to \mathcal{L}(G)$. Its inverse is (immediately) given by the map

$$\mathcal{L}(G) \to T_e G, \qquad X \mapsto X(e) \in T_e G.$$

2.5 $T_e G \cong \mathcal{L}(G)$ as Lie Algebras

To show that $T_e G$ and $\mathcal{L}(G)$ are isomorphic as Lie algebras as well as vector fields, we must show that the map

$$\nu \colon T_e G \to \mathcal{L}(G)$$

preserves the brackets, i.e.

$$\nu_{\mathrm{ad}(\xi)\eta} = \left[\nu_{\xi}, \nu_{\eta}\right]$$

for all $\xi, \eta \in T_e G$. Since the Lie bracket of vector fields can be described easily in terms of flows, it might be helpful to know what the flows of these vector fields look like.

Claim 2.4. Let $\xi \in T_eG$. and $g \in G$. Then the flow of ν_{ξ} through g is the curve $c \colon \mathbb{R} \to G$ given by

$$c(t) = L_g \circ \exp(t\xi).$$

Proof. Note that $c(0) = L_g \circ \exp(\vec{0}) = L_g(e) = g$. Let $t \in \mathbb{R}$. Then

$$\begin{aligned} c(t) &= \frac{d}{ds} \big|_{s=t} c(t) = \frac{d}{ds} \big|_{s=0} c(s+t) \\ &= \frac{d}{ds} \big|_{s=0} L_g \circ \exp\left((s+t)\xi\right) = \frac{d}{ds} \big|_{s=0} L_g \circ \exp\left((t+s)\xi\right) \\ &= \frac{d}{ds} \big|_{s=0} L_g \left(\exp(t\xi) \cdot \exp(s\xi)\right) = \frac{d}{ds} \big|_{s=0} L_g \circ L_{\exp(t\xi)} \left(\exp(s\xi)\right) \\ &= \frac{d}{ds} \big|_{s=0} L_g \exp(t\xi) \left(\exp(s\xi)\right) \\ &= \left(T_e L_g \exp(t\xi)\right) \left(\frac{d}{ds} \big|_{s=0} \exp(s\xi)\right) \\ &= \left(T_e L_g \exp(t\xi)\right) (\xi) \\ &= \nu_\xi \left(g \exp(t\xi)\right) \\ &= \nu_\xi (c(t)). \end{aligned}$$

QED

Theorem 2.5. Let $\xi, \eta \in T_eG$. Then

$$\nu_{ad(\xi)\eta} = \left[\nu_{\xi}, \nu_{\eta}\right]$$

Proof. Recall that the flow of ν_{ξ} at time $t \in \mathbb{R}$ is the map $G \to G$ given by $R_{\exp(t\xi)}$. Let $g \in G$. Then using the definition of Ψ , Ad, and ν , the linearity of tangent maps, and Proposition ??, we calculate

$$\begin{split} [\nu_{\xi}, \nu_{\eta}](g) &= \frac{d}{dt} \big|_{t=0} \Big(\big(R_{\exp(t\xi)} \big) \nu_{eta} \Big)(g) \\ &= \frac{d}{dt} \big|_{t=0} T R_{\exp(t\xi)} \circ \nu_{eta} \circ R_{\exp(t\xi)}^{-1}(g) \\ &= \frac{d}{dt} \big|_{t=0} T R_{\exp(t\xi)} \circ \nu_{eta} \big(g \exp(-t\xi) \big) \\ &= \frac{d}{dt} \big|_{t=0} T R_{\exp(t\xi)} \circ T L_{g \exp(-t\xi)}(\eta) \end{split}$$

$$\begin{split} &= \frac{d}{dt} \big|_{t=0} T \Big(R_{\exp(t\xi)} \circ L_{g \exp(-t\xi)} \Big) (\eta) \\ &= \frac{d}{dt} \big|_{t=0} T \Big(R_{\exp(t\xi)} \circ L_{g} \circ L_{\exp(-t\xi)} \Big) (\eta) \\ &= \frac{d}{dt} \big|_{t=0} T \Big(L_{g} \circ R_{\exp(t\xi)} \circ L_{(\exp(t\xi))^{-1}} \Big) (\eta) \\ &= \frac{d}{dt} \big|_{t=0} \big(TL_{g} \big) \circ \big(T\Psi_{\exp(t\xi)} \big) (\eta) \\ &= \frac{d}{dt} \big|_{t=0} \big(TL_{g} \big) \circ \operatorname{Ad}(\exp t\xi) (\eta) \\ &= \big(TL_{g} \big) \Big(\frac{d}{dt} \big|_{t=0} \operatorname{Ad}(\exp t\xi) \eta \Big) \\ &= \big(TL_{g} \big) [\xi, \eta] \\ &= \nu_{[\xi,\eta]}(g). \end{split}$$

QED