Limited choice and randomness in evolution of networks Lecture 1

Cornell Probability summer school July 2012

Shankar Bhamidi

Department of Statistics and Operations Research University of North Carolina

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Plan of the lectures

Underlying theme

- Mathematical techniques for dynamic random graph models
- Effect of limited choice

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- Mathematical techniques for dynamic random graph models
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Lecture content

- Lecture 1: Critical random graphs, Bounded size rules [Scaling limits]
- Lecture 2: Preferential attachment models, random trees [Local Weak convergence]

Lecture 1

- General motivation
- Critical random graphs
- Bounded size rules and the emergence of the giant
- Method of proof (Joint work with Budhiraja and Wang)
- Extensions: "Explosive percolation"

Lecture 2

- Preferential attachment models
- Convergence of random trees
- Implications: Convergence of the spectral measure (Arnab Sen, Steve Evans, SB)
- Power of choice in random trees
- Local weak convergence (Angel, Pemantle, SB)

Application setting

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- Limited choice Choose 2 bins u.a.r.
- Put ball in bin with minimal # of balls at that stage

Max load $\sim \Theta(\log \log n)$

Network models

Motivation

- Last few years have seen an explosion in empirical data on real world networks.
- Has motivated an interdisciplinary study in understanding the emergence of properties of these network models.
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Limited choice

- Incorporate effect of limited choice in network formation
- Mathematically understand explosive percolation
- Simple variants of standard models give much better fit but hard to mathematically analyze



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Beautiful math theory

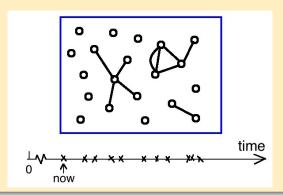


Bounded size rules

- $G_n(0) = \mathbf{0}_n$ the graph with n vertices but no edges
- Each step, choose one edge e uniformly among all $\binom{n}{2}$ possible edges, and add it to the graph.

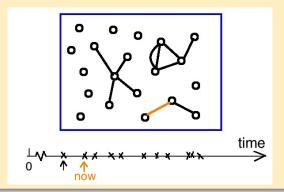
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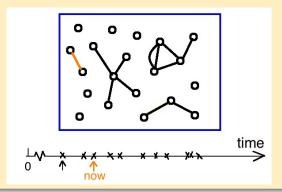
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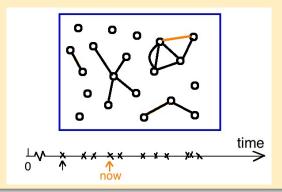


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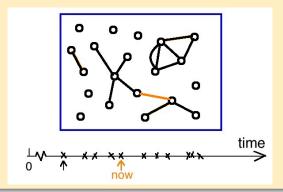
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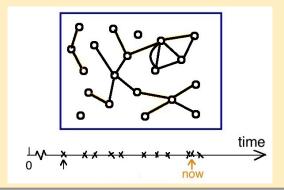
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- (sub-critical) when t < 1, $C_n^{(1)} = O(\log n)$, $C_n^{(2)} = O(\log n)$.
- (critical) when t=1, $\mathcal{C}_n^{(1)}\sim n^{2/3}$, $\mathcal{C}_n^{(2)}\sim n^{2/3}$.

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- after initial work by [ER1960], further work by [JKLP1994], finally proved by [Aldous1997].
- Merging dynamics through the scaling window of the components described by a Markov Process called the multiplicative coalescent.
- Formal existence of multiplicative coalescent.

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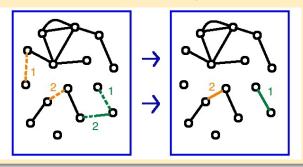
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The Bohman-Frieze process

[Bohman, Frieze 2001] The delay of phase transition

Consider the continuous time version $\mathcal{G}_n^{\rm BF}(t)$, then there exists $\epsilon>0$ such that at time $t_c^{\rm ER}+\epsilon$,

$$C_n^{(1)}(t_c^{ER} + \epsilon) = o(n)$$

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Near Criticality

- Janson and Spencer (2011) analyzed how $s_2(\cdot), s_3(\cdot) \to \infty$ as $t \uparrow t_c$.
- Kang, Perkins and Spencer (2011) analyze the near subcritical $(t_c \epsilon)$ regime.

General bounded size rules

- Fix K > 1
- Let $\Omega_K = \{1, 2, \dots, K, \omega\}$
- General bounded size rule: subset $F \subset \Omega^4_K$.
- Pick 4 vertices uniformly at random. If $(c(v_1),c(v_2),c(v_3),c(v_4)) \in F$ then choose edge e_1 else e_2

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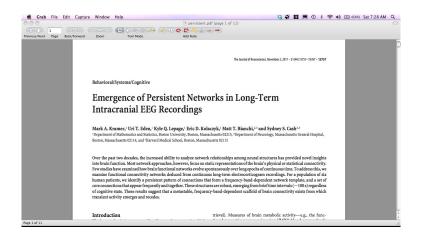
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BF model

$$K = 1, F = \{(1, 1, \alpha, \beta)\}.$$



Applied context



Main questions

• Question: when $t=t_c$, do we have $\mathcal{C}_n^{(1)}\sim n^{2/3}$? How do components merge?

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- What about the surplus of the largest components in the scaling window?

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- $l_{\downarrow}^2 = \{(x_i)_{i \ge 1} : x_1 \ge x_2 \ge \dots \ge 0, \sum_i x_i^2 < \infty \}$
- $l_{\downarrow}^{2,*} = \left\{ (x_i, y_i)_{i \ge 1} : (x_i) \in l_{\downarrow}^2, y_i \in \mathbb{Z}_+, \sum_i x_i y_i < \infty \right\}$
- $d((x,y),(x',y')) = \sqrt{\sum_i (x_i x_i')^2} + \sum_i |x_i y_i x_i' y_i'| + \sum_{i=1}^{\infty} \frac{|y_i y_i'|}{2^i}$



The Erdős-Rényi random graph

Theorem (Aldous 1997)

Let $(\mathcal{C}_n^{(1)}(t), \mathcal{C}_n^{(2)}(t), ...)$ be the component sizes of $\mathcal{G}_n^{ER}(t)$ in decreasing order and $\xi_i(t)$ the corresponding complexity (surplus). Define rescaled size vector $\mathbf{C}_n^*(\lambda), -\infty < \lambda < +\infty$ as

$$\left(\left(\frac{1}{n^{2/3}}\mathcal{C}_n^{(i)}\left(t_c + \frac{\lambda}{n^{1/3}}\right),\right.\right)$$

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Then $\mathbf{C}_n(\lambda) \stackrel{d}{\longrightarrow} \mathbf{X}(\lambda) = (X(\lambda), \xi(\lambda))$. Here $(X(\lambda), -\infty < \lambda < +\infty)$ is the standard multiplicative coalescent, a continuous time Markov process on the state space l_{\perp}^2 .

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Distribution for fixed λ

• For fixed $\lambda \in \mathbb{R}$, let

$$W_{\lambda}(t) = W(t) + \lambda t - \frac{t^2}{2}$$

- $W_{\lambda}(t) = W(t) + \lambda t \frac{t^2}{2},$ $\bar{W}_{\lambda}(\cdot)$ is the above process reflected at 0.
- $X(\lambda)$ has same distribution as lengths of excursions away from 0 of $\bar{W}(\cdot)$ arranged in decreasing order

- ullet Think of each edge having a Poisson rate 1/n clock
- $\bar{C}_i(\lambda) = n^{-2/3}C_i(1 + \lambda/n^{1/3})$
- At some time λ , rate at which two components i, j merge:

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$$C_i(1 + \frac{\lambda}{n^{1/3}})C_j(1 + \frac{\lambda}{n^{1/3}}) \cdot \frac{1}{n} \cdot \frac{1}{n^{1/3}} = \bar{C}_i(\lambda)\bar{C}_j(\lambda)$$

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Dynamics of $\mathbf{X}(\lambda)$

- suppose $\mathbf{X}(\lambda) = (x_1, x_2, x_3, ...)$, each x_l is viewed as the size of a cluster.
- each pair of clusters of sizes (x_i,x_j) merges at rate x_ix_j into a cluster of size x_i+x_j .
- if x_i, x_j is merging, then $(x_1, x_2, x_3, ...) \rightsquigarrow (x'_1, x'_2, x'_3, ...)$ where the latter is the re-ordering of $\{x_i + x_j, x_l : l \neq i, j\}$.

Bounded size rules

Theorem (Bhamidi, Budhiraja, Wang, 2012)

Let $(\mathcal{C}_n^{(1)}(t), \mathcal{C}_n^{(2)}(t), ...)$ be the component sizes of $\mathcal{G}_n^{BSR}(t)$ in decreasing order and $\xi_i(t)$ the corresponding surplus. Define the rescaled size vector $\mathbf{C}_n(\lambda)$, $-\infty < \lambda < +\infty$ as the vector

$$((\bar{C}_i(\lambda), \xi_i(\lambda) : i \ge 1) = \left(\frac{\beta^{1/3}}{n^{2/3}} C_n^{(i)} (t_c + \frac{\beta^{2/3} \alpha \lambda}{n^{1/3}}), \xi_i (t_c + \frac{\beta^{2/3} \alpha \lambda}{n^{1/3}}) : i \ge 1\right)$$

where α, β are constants determined by the BSR process. Then

$$\mathbf{C}_n(\lambda) \stackrel{d}{\longrightarrow} \mathbf{X}(\lambda)$$

where $(\mathbf{X}(\lambda), -\infty < \lambda < +\infty)$ is the standard augmented multiplicative coalescent and convergence happens in l^2_\downarrow with metric d.



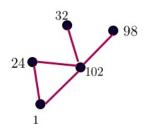
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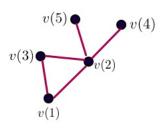
- Branching process methods: Great tool above and below criticality.
- Exploration walks: Very refined results presence of lots of independence, including structure of components
- **Differential equation method:** Technical, standard workhorse for such models. Can be pushed all the way to the critical window.

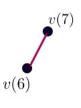
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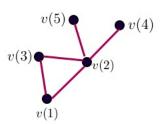
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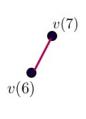
$$c(1) = 2$$

$$c(2) = 2$$

$$c(3) = 0$$

. . .





Typical method of proof

Exploration of the graph

- Explore the components of the graph one by one
- ullet choose a vertex. Let c(1) be the number of children of this vertex
- ullet choose one of the children of this vertex, let c(2) be number of children of this vertex
- continue, when one component completed move onto another component
- Define Z(0) = 0, Z(i) = Z(i-1) + c(i) 1
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Bounded size rules

Hard to think about exploration process especially at criticality

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- Hard to think about exploration process especially at criticality
- Turns out: Easier to analyze the entire process

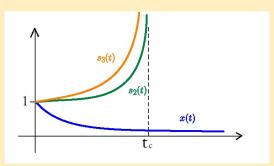
Proof idea: The Bohman-Frieze process

Where does t_c come from ?

Define $X_n(t) = \#$ of singletons, $S_2(t) = \sum_i (\mathcal{C}_n^{(i)}(t))^2$, $S_3(t) = \sum_i (\mathcal{C}_n^{(i)})^3$. and $\bar{x}_n(t) = X_n(t)/n$, $\bar{s}_2(t) = S_2/n$, $\bar{s}_3(t) = S_3/n$.

Then [Spencer, Wormald 2004] for any fix
$$t > 0$$
,

$$\bar{x}_n(t) \stackrel{\mathbb{P}}{\longrightarrow} x(t), \qquad \bar{s}_2(t) \stackrel{\mathbb{P}}{\longrightarrow} s_2(t), \qquad \bar{s}_3(t) \stackrel{\mathbb{P}}{\longrightarrow} s_3(t)$$



Behavior of $x_n(t)$

• In small time interval $[t, t + \Delta(t))$, $x_n(t) \to x_n(t) - 1/n$ at rate

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• Similar analysis suggests that for $\bar{s}_2(t), \bar{s}_3(t)$

$$\begin{split} s_2'(t) &= x^2(t) + (1 - x^2(t))s_2^2(t) & \text{for } t \in [0, t_c), \qquad s_2(0) = 1 \\ s_3'(t) &= 3x^2(t) + 3(1 - x^2(t))s_2(t)s_3(t) & \text{for } t \in [0, t_c), \qquad s_3(0) = 1. \end{split}$$

The Bohman-Frieze process

Scaling exponents of s_2 and s_3 (Janson, Spencer 11)

- Functions $x(t), s_2(t), s_3(t)$ are determined by some differential equations
- Differential equations imply \exists constants α, β such that $t \uparrow t_c$

$$s_2(t) \sim \frac{\alpha}{t_c - t}$$

$$s_3(t) \sim \beta(s_2(t))^3 \sim \beta \frac{\alpha^3}{(t_c - t)^3}$$

I: Regularity conditions of the component sizes at

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$$-\infty$$
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• Let
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- For $\delta \in (1/6,1/5)$ let $t_n=t_c-n^{-\delta}=t_c+\beta^{2/3}\alpha\frac{\lambda_n}{n^{1/3}}$, then $\lambda_n=-\beta^{2/3}\alpha n^{1/3-\delta}$.

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- Need to verify the three conditions

$$\frac{\sum_{i} \left(\bar{C}_{i}(\lambda_{n})\right)^{3}}{\left[\sum_{i} \left(\bar{C}_{i}(\lambda_{n})\right)^{2}\right]^{3}} \xrightarrow{\mathbb{P}} 1 \qquad \Leftrightarrow \frac{n^{2}S_{3}(t_{n})}{S_{2}^{3}(t_{n})} \xrightarrow{\mathbb{P}} \beta$$

$$\frac{1}{\sum_{i} \left(\bar{C}_{i}(\lambda_{n})\right)^{2}} + \lambda_{n} \xrightarrow{\mathbb{P}} 0 \qquad \Leftrightarrow \frac{n^{4/3}}{S_{2}(t_{n})} - \frac{n^{-\delta+1/3}}{\alpha} \xrightarrow{\mathbb{P}} 0$$

$$\frac{\bar{C}_{1}(\lambda_{n})}{\sum_{i} \left(\bar{C}_{i}(\lambda_{n})\right)^{2}} \xrightarrow{\mathbb{P}} 0 \qquad \Leftrightarrow \frac{n^{2/3}C_{n}^{(1)}(t_{n})}{S_{2}(t_{n})} \xrightarrow{\mathbb{P}} 0$$

II: Dynamics of merging in the critical window

The dynamic of merging

 \bullet In any small time interval [t,t+dt), two components i and j merge at rate

$$\frac{2}{n^3} \left[\binom{n}{2} - \binom{X_n(t)}{2} \right] \mathcal{C}_i(t) \mathcal{C}_j(t)$$

$$\sim \frac{1}{n} (1 - \bar{x}^2(t)) \mathcal{C}_i(t) \mathcal{C}_j(t)$$

Let $\lambda=(t-t_c)n^{1/3}/\alpha\beta^{2/3}$ be rescaled time paramter, rate at which two components merge

$$\gamma_{ij}(\lambda) \sim \frac{(1 - x^2(t_c + \beta^{2/3}\alpha \frac{\lambda}{n^{1/3}}))}{n} \frac{\beta^{2/3}\alpha}{n^{1/3}} C_i \left(t_c + \frac{\beta^{2/3}\alpha\lambda}{n^{1/3}}\right) C_j \left(t_c + \frac{\beta^{2/3}\alpha\lambda}{n^{1/3}}\right)$$

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$$\begin{split} \gamma_{ij}(\lambda) &\sim \frac{(1-x^2(t_c+\beta^{2/3}\alpha\frac{\lambda}{n^{1/3}}))}{n} \frac{\beta^{2/3}\alpha}{n^{1/3}} \mathcal{C}_i \left(t_c+\frac{\beta^{2/3}\alpha\lambda}{n^{1/3}}\right) \mathcal{C}_j \left(t_c+\frac{\beta^{2/3}\alpha\lambda}{n^{1/3}}\right) \\ &= \alpha \left(1-x^2\left(t_c+\beta^{2/3}\alpha\frac{\lambda}{n^{1/3}}\right)\right) \bar{\mathcal{C}}_i(\lambda) \bar{\mathcal{C}}_j(\lambda) \\ &= \bar{\mathcal{C}}_i(\lambda) \bar{\mathcal{C}}_j(\lambda) \qquad \text{since } \alpha(1-x^2(t_c)) = 1 \end{split}$$

How to check regularity conditions

Analysis of $\mathcal{C}_n^{\scriptscriptstyle (1)}(t)$

- Key point: need to get refined bounds on maximal component in barely subcritical regime.
- known result: for fixed $t < t_c$, $C_n^{\text{\tiny (1)}}(t) = O(\log n)$.
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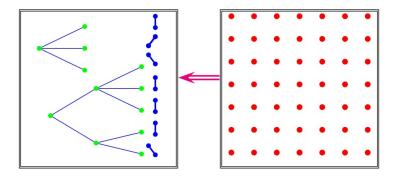
Lemma (Bounds on the largest component)

Let $\delta \in (0,1/5)$, t_c be the critical time for the BF process, $\mathcal{C}_n^{\scriptscriptstyle (1)}(t)$ be the size of the largest component. Then there exists a constant $B=B(\delta)$ such that as $n\to +\infty$,

$$\mathbb{P}\{\mathcal{C}_n^{\scriptscriptstyle (1)}(t) \leq \frac{B\log^4 n}{(t_c-t)^2} \text{ for all } t < t_c - n^{-\delta}\} \to 1$$

Proof strategy: Coupling with a near critical multi-type branching process on an infinite dimensional type space. delicate analysis of the maximal eigenvalue.

Random graph with Immigrating doubletons



Sketch of the proof

Regularity condition at time $\lambda = -\infty$

Check the following properties for the un-scaled component sizes. For $\delta \in (1/6, 1/5)$, and $t_n = t_n - n^{-\delta}$.

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Analysis of $S_2(t)$, $S_3(t)$

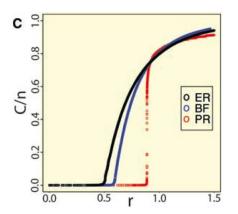
- above relations hold for limiting functions s_2, s_3 for t_n .
- Delicate stochastic analytic argument combined with result on $\mathcal{C}_n^{(1)}(t)$ to show this holds for S_2, S_3 near criticality.



Work in progress

Explosive percolation

In 2009, Achlioptas, D'Souza and Spencer considered "product rule". Conjectured that this process exhibits Explosive percolation



Fix K

- ullet Choose 2 edges $e_1=(v_1,v_2)$ and $e_2=(v_3,v_4)$ at random
- If $\max\{C(v_1),C(v_2)C(v_3),C(v_4)\} \leq K$, then use the edge which minimizes of $\min\{C(v_1)C(v_2),C(v_3)C(v_4)\}$.
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Consider the rescaled and re-centered component sizes

$$\mathbf{C}_K(\lambda) = \left(\frac{1}{n^{2/3}} \mathcal{C}_K^{(i)} \left(t_c(K) + \gamma(K) \frac{\lambda}{n^{1/3}} \right) : i \ge 1 \right) \qquad \lambda \in \mathbb{R}$$

Then we have $(\mathbf{C}_K(\lambda):\lambda\in\mathbb{R})\stackrel{d}{\longrightarrow} (X(\lambda):\lambda\in\mathbb{R})$ as $n\to\infty$.

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Other questions

"Natural questions"

- What happens if we start with a configuration other than the empty graph?
- Related to the entrance boundary of the multiplicative coalescent.

Unnatural next questions

Scaling limits?

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Unnatural next questions

- Scaling limits?
- Conjecture: Rescale each edge by $n^{-1/3}$
- Largest components converge to random fractals (Gromov-Hausdorff sense), the same limits as for Erdos-Renyii

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- Enormous literature in applied sciences
- Deep connections to statistical physics models of disorder

Statistical physics models of disorder

Weak disorder (First passage percolation)

- Weight of a path = sum of weight on edges
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Strong disorder (Minimal spanning tree)

- Weight of path = max edge on the path
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Adario-Berry, Broutin, Goldschmidt + Miermont

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- This implies that $n^{-1/3}\mathcal{M}_n$ converges to a limiting random fractal [BBGM]
- Open Problem: Show that for these models, have same limiting structure

